

Exact expressions of mean first-passage times and splitting probabilities for random walks in bounded rectangular domains

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Recently, we have proposed a novel computation method of first-passage times of a random walker between a starting site and a target site of regular bounded lattices of arbitrary shape.¹ Such first passage time properties are, for example, crucial in describing the kinetics of diffusion limited reactions in confined media.^{2,3} The obtained expressions involve pseudo-Green functions,⁴ which have been estimated according to different approximation schemes. In this Note, we give the exact expression of these pseudo-Green functions in both cases of a rectangular domain with reflecting boundaries and a rectangular domain with periodic boundary conditions. This allows us to provide exact and explicit expressions of mean first-passage times with one or two target sites.

Consider first a random walker starting at the source S of position \mathbf{r}_S of a rectangular regular lattice with reflecting boundary conditions. The mean time it takes to reach the target T of position \mathbf{r}_T for the first time is given by¹

$$\langle \mathbf{T} \rangle = N[H(\mathbf{r}_T|\mathbf{r}_T) - H(\mathbf{r}_T|\mathbf{r}_S)], \quad (1)$$

where H is the pseudo-Green function, which satisfies

$$H(\mathbf{r}_i|\mathbf{r}_j) = \frac{1}{\sigma} \sum_{k \in N_i} H(\mathbf{r}_k|\mathbf{r}_j) + \delta_{ij} - \frac{1}{N}, \quad (2)$$

σ being the coordination number of the lattice, namely, 4 for a two-dimensional (2D) lattice or 6 for a three-dimensional (3D) lattice. In this expression, N_i is the set of neighbors of

the site i , considering that a site which is at the boundary of the domain is its own neighbor. N is the total number of sites of the lattice. In the presence of two absorbing targets, the mean time it takes to reach either of the two targets is

$$\langle \mathbf{T} \rangle = N \frac{(H_{01} - H_{1s})(H_{02} - H_{2s}) - (H_{12} - H_{2s})(H_{12} - H_{1s})}{H_{01} + H_{02} - 2H_{12}}, \quad (3)$$

while the eventual hitting probabilities P_i to reach the target i writes

$$P_1 = \frac{H_{1s} + H_{02} - H_{2s} - H_{12}}{H_{01} + H_{02} - 2H_{12}}, \quad (4)$$

$$P_2 = \frac{H_{2s} + H_{01} - H_{1s} - H_{12}}{H_{01} + H_{02} - 2H_{12}},$$

where $H_{12} = H(\mathbf{r}_{T_1}|\mathbf{r}_{T_2})$ and, for $i = 1$ or 2 , $H_{is} = H(\mathbf{r}_{T_i}|\mathbf{r}_S)$ and $H_{0i} = H(\mathbf{r}_{T_i}|\mathbf{r}_{T_i})$.

The exact expression of the pseudo-Green function H involved in previous equations may be computed explicitly with the help of Fourier analysis. For a 2D domain with X sites in the x direction and Y sites in the y direction, H writes

$$H(\mathbf{r}|\mathbf{r}') = \frac{1}{N} \sum_{m=1}^{X-1} \sum_{n=1}^{Y-1} \frac{4 \cos(n\pi x'/X) \cos(n\pi y'/Y) \cos(m\pi x/X) \cos(n\pi y/Y)}{1 - (1/2)(\cos(m\pi/X) + \cos(n\pi/Y))} \\ + \frac{1}{N} \sum_{m=1}^{X-1} \frac{4 \cos(m\pi x'/X) \cos(m\pi x/X)}{1 - \cos(m\pi/X)} + \frac{1}{N} \sum_{n=1}^{Y-1} \frac{4 \cos(n\pi y'/Y) \cos(n\pi y/Y)}{1 - \cos(n\pi/Y)}, \quad (5)$$

where x and y are the coordinates of \mathbf{r} and x' and y' those of \mathbf{r}' . Here, the coordinates of the left-bottom corner are taken to be equal to $(1/2, 1/2)$, so that all the coordinates are half-integers.

Figure 1 shows the mean absorption time $\langle \mathbf{T} \rangle$ given by Eq. (3) as a function of the target positions together with

numerical simulations in the case of a 2D square domain. These simulations have been performed with a method based on the exact enumeration method.⁵

Similar results can be obtained in the case of a rectangular domain with periodic boundary conditions. Here, the pseudo-Green function is

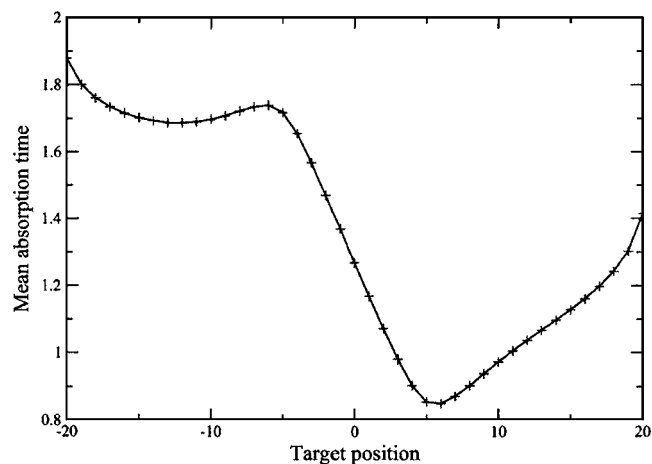


FIG. 1. 2D two-target simulations (red crosses) vs theory (plain line). One target is fixed at $(-5,0)$; the source is fixed at $(5,0)$; the other target is at $(x,3)$. The domain is a square of side 41; the middle is the point $(0,0)$. The absorption time is normalized by the number of sites N .

$$H(\mathbf{r}|\mathbf{r}') = \frac{1}{N} \sum_{m=0}^{X-1} \sum_{n=\delta_{0m}}^{Y-1} \frac{\exp[2im\pi(x-x')/X] \exp[2ni\pi(y-y')/Y]}{1 - (1/2)(\cos(2m\pi/X) + \cos(2n\pi/Y))}, \quad (6)$$

where δ_{0m} is the Kronecker delta.

Note that our results can be easily extended to the 3D case.

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¹S. Condamin, O. Bénichou, and M. Moreau, Phys. Rev. Lett. **95**, 260601 (2005).

²S. Rice, *Diffusion-Limited Reactions* (Elsevier, Amsterdam, 1985).

³S. Redner, *A guide to First-Passage Processes* (Cambridge University Press, Cambridge, 2001).

⁴G. Barton, *Elements of Green's Functions and Propagation* (Oxford Science, Oxford, 1989).

⁵I. Majid, D. B. Avraham, S. Havlin, and H. E. Stanley, Phys. Rev. B **30**, 1626 (1984).